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ADAPTIVE GDSW COARSE SPACES FOR OVERLAPPING SCHWARZ METHODS IN THREE DIMENSIONS

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Abstract. A robust two-level overlapping Schwarz method for scalar elliptic model problems with highly varying coefficient functions is introduced. While the convergence of standard coarse spaces may depend strongly on the contrast of the coefficient function, the condition number bound of the new method is independent of the coefficient function. Its coarse space is based on discrete harmonic extensions of vertex, edge, and face interface functions, which are computed from the solutions of corresponding local generalized edge and face eigenvalue problems. The local eigenvalue problems are of the size of the edges and faces of the decomposition, and the eigenvalue problems can be constructed solely from the local subdomain stiffness matrices and the fully assembled global stiffness matrix. The new AGDSW (Adaptive Generalized Dryja–Smith–Widlund) coarse space always contains the classical GDSW coarse space by construction of the generalized eigenvalue problems. Numerical results supporting the theory are presented for several model problems in three dimensions using structured as well as unstructured meshes and unstructured decompositions.

Key words. domain decomposition, multiscale, GDSW, overlapping Schwarz, adaptive coarse spaces

AMS subject classifications. 65F08,65F10,65N55,68W10

1. Introduction. We introduce an adaptive coarse space for the two-level overlapping Schwarz method [50, 53] and prove a condition number bound that is independent of heterogeneities in the coefficient function of the underlying variational problem (1); cf. section 6 and the supporting numerical results in section 9. The presented coarse space – adaptive GDSW (AGDSW) – can be regarded as an extension of the energy-minimizing GDSW coarse space (Generalized Dryja–Smith–Widlund) [7, 6], as the latter of which is always contained in the former space.

The classical GDSW coarse space is constructed by an energy-minimal extension of null space functions on the interface such that the kernel of the elliptic operator is represented. This can also be carried out algebraically and results in a method that is robust for a class of coefficient functions; cf., e.g., [6, Table 5.3] and [21, Chapter 5].

The GDSW method has been applied to a variety of model problems, see, e.g., [8, 9] for the application to linear elasticity. In [26], the use of GDSW was applied to the highly nonlinear structural part in fluid-structure interaction simulations, and in [22], it was applied to various saddle point problems. A parallel implementation of GDSW is publicly available as the FROSch [25] preconditioner (Fast and Robust Overlapping Schwarz) as part of the Trilinos [31] package ShyLU [48]; for implementation details and numerical results, see [26, 28, 27]. Furthermore, recently, in [29], a three-level parallel implementation of GDSW in two dimensions was presented. Reduced dimension GDSW coarse spaces have been considered, e.g., in [9, 5]; see also [11] and the references therein, and [30] for results on the parallel performance.

However, classical GDSW coarse spaces are not sufficient to obtain a method which is robust for arbitrary coefficient jumps; see, e.g., [21, Chapter 5]. To this end, adaptive (w.r.t. the coefficient function) coarse spaces have been developed in

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the field of domain decomposition methods.

A natural initial choice for basis functions to treat coefficient variations are multiscale finite element (MsFEM) functions [33, 15]; see [1, 4, 19, 24, 16]. To define MsFEM functions, boundary values need to be chosen carefully which in [19, 24, 16] requires solving a problem on the interface. By contrast, the construction of GDSW vertex-based functions, which are included in the AGDSW space, is much simpler.

In addition to vertex-based functions, the construction of the AGDSW coarse space uses the energy-minimal extensions of low-frequency eigenmodes on the edges and faces of the domain decomposition. Note that our approach is different from the two-dimensional AGDSW coarse space in [23] which allows for a simpler implementation and can lead to smaller coarse spaces according to our numerical experiments. A special emphasis is placed on the reduction of the coarse space dimension by also integrating energy-minimal extensions into the eigenvalue problems; cf. the numerical results in section 9. This strategy has also been used for the coarse space in [24], which was inspired by a special finite element method based on approximate component mode synthesis [32].

Local generalized eigenvalue problems to construct coarse spaces have been used earlier to obtain methods which are robust to coefficient jumps. In [16], the authors present two approaches to construct coarse spaces which lead to eigenvalue problems of the same size as here. The setup of their eigenvalue problems is cheaper, however, the coarse space dimension can be significantly larger; cf. section 9. In [17], Galvis and Efendiev use generalized eigenvalue problems on unions of subdomains resulting in large eigenvalue problems. Dolean et al. proposed generalized eigenvalue problems on subdomain boundaries but a restriction on the class of coefficient functions was required to prove the condition number bound in [12]. In [51], Spillane et al. then introduced the coarse space GenEO in which they reduced the generalized eigenvalue problems to the overlap of subdomains allowing arbitrary coefficient functions. A further reduction to edges in two dimensions, and edges and faces in three dimensions, was realized in [18, 19, 24, 23, 16]. Other notable contributions to multiscale domain decomposition for overlapping Schwarz are, e.g., [20, 3]. Adaptive coarse spaces for nonoverlapping domain decomposition methods have gained much interest as well; see, e.g., [2, 43, 44, 52, 46, 36, 38, 47, 35, 37, 45].

2. Model problem. On a polyhedral domain $\Omega \subset \mathbb{R}^3$, we consider the variational problem: find $u \in H_0^1(\Omega)$, such that

$$(1) \quad a_\Omega(u, v) = L(v) \quad \forall v \in H_0^1(\Omega),$$

$$\text{where } a_\Omega(u, v) := \int_\Omega A(x)(\nabla u(x))^T \nabla v(x) dx \quad \text{and} \quad L(v) := \int_\Omega f(x)v(x) dx,$$

respectively, and where $A: \mathbb{R}^3 \rightarrow \mathbb{R}$ is a scalar coefficient function and $f \in L^2(\Omega)$. In this paper, A is typically highly heterogeneous, possibly having discrete values with large variations. In addition to that, we denote the semi-norm corresponding to the bilinear form $a_\Omega(\cdot, \cdot)$ as

$$|u|_{a, \Omega}^2 := a_\Omega(u, u).$$

We assume that the coefficient function $A(x)$ satisfies

$$0 < A_{\min} \leq A(x) \leq A_{\max} \quad \forall x \in \bar{\Omega}.$$

Then, the Lax-Milgram lemma guarantees a unique solution of (1). Let

$$(2) \quad Ku = f$$

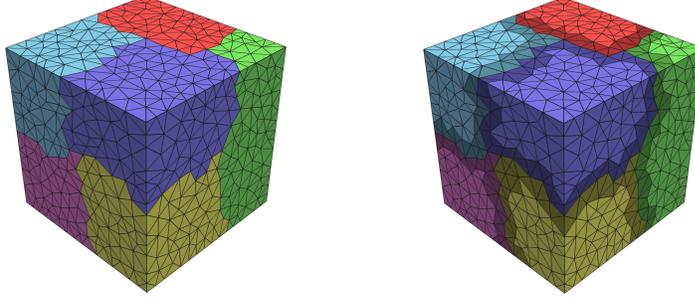


FIG. 1. Example of a nonoverlapping domain decomposition $\{\Omega_i\}_{i=1}^N$ (left) and the corresponding overlapping domain decomposition $\{\Omega'_i\}_{i=1}^N$ (right) with overlap $\delta = 1h$.

be the discretization of problem (1) by piecewise linear or trilinear finite elements on a triangulation τ_h . Here, K is the stiffness matrix, f the right hand side, and u the vector corresponding to the finite element solution in the finite element space $V^h(\Omega)$. Throughout this paper, we assume that the coefficient function A is constant on each finite element $T \in \tau_h$. However, our method is not restricted to these cases. In order to solve this problem, we use the conjugate gradient method preconditioned by a two-level overlapping Schwarz preconditioner.

3. Two-level overlapping Schwarz methods. Let $\{\Omega_i\}_{i=1}^N$ be a nonoverlapping domain decomposition of Ω into polyhedral subdomains Ω_i with a typical subdomain diameter of H . The interface Γ of the nonoverlapping domain decomposition is defined as $\Gamma = \bigcup_{i=1}^N \partial\Omega_i \setminus \partial\Omega$.

Next, we obtain a corresponding overlapping decomposition $\{\Omega'_i\}_{i=1}^N$ of Ω by extending the nonoverlapping subdomains by k layers of finite elements. This results in an overlap $\delta = kh$; cf. Figure 1. We define as $R_i : V^h(\Omega) \rightarrow V_i := V^h(\Omega'_i)$, $i = 1, \dots, N$, the restriction to the local finite element space on the overlapping subdomain Ω'_i ; R_i^T is the corresponding prolongation to $V^h(\Omega)$. In addition, let V_0 be some global coarse space and $R_0 : V^h(\Omega) \rightarrow V_0 \subset V^h(\Omega)$ the corresponding coarse interpolation. We use exact solvers, and therefore the local and coarse bilinear forms on the subspaces are given by

$$\tilde{a}_i(u_i, v_i) = a_\Omega(R_i^T u_i, R_i^T v_i) \quad \forall u_i, v_i \in V_i,$$

$i = 0, \dots, N$. Then, the additive two-level Schwarz operator is given by

$$P_{\text{OS-2}} = M_{\text{OS-2}}^{-1} K = R_0^T K_0^{-1} R_0 + \sum_{i=1}^N R_i^T K_i^{-1} R_i K$$

with local stiffness matrices $K_i = R_i K R_i^T$, for $i = 1, \dots, N$; cf. [53, Chapter 2.2] and coarse operator $K_0 = R_0 K R_0^T$.

The condition number of the two-level Schwarz operator for the finite element problem (2) using Lagrangian coarse basis functions for K_0 depends on the contrast of the coefficient function A , i.e.,

$$\kappa(M_{\text{OS-2}}^{-1} K) \leq C \max_{T \in \tau_H} \max_{x, y \in \omega_T} \left(\frac{A(x)}{A(y)} \right) \left(1 + \frac{H}{\delta} \right);$$

cf. [20]. Here, τ_H corresponds to the set of all coarse mesh elements, and ω_T to the union of all coarse mesh elements which touch a coarse mesh element T . This bound can be improved but a dependence on the coefficient contrast remains.

4. The GDSW preconditioner. The GDSW preconditioner [6, 7] is a two-level additive overlapping Schwarz preconditioner with exact solvers as described in the previous section. Thus, the preconditioner can be written in the form

$$M_{\text{GDSW}}^{-1} = \Phi K_0^{-1} \Phi^T + \sum_{i=1}^N R_i^T K_i^{-1} R_i,$$

where $\Phi = R_0^T$. Here, the columns of Φ are coefficient vectors corresponding to the coarse basis functions and the main ingredient of the GDSW preconditioner.

The interface can be decomposed as $\Gamma = \left(\bigcup_{f \in \mathcal{F}} f\right) \cup \left(\bigcup_{e \in \mathcal{E}} e\right) \cup \left(\bigcup_{v \in \mathcal{V}} v\right)$, where \mathcal{F} is the set of all faces, \mathcal{E} the set of all edges, and \mathcal{V} the set of all vertices; see, e.g., [41, Sect. 3] and [39, 40, Sect. 2]. The discrete characteristic functions χ_*^h of the vertices, edges, and faces form a partition of unity on Γ , i.e.,

$$1 = \sum_{v \in \mathcal{V}} \chi_v^h + \sum_{e \in \mathcal{E}} \chi_e^h + \sum_{f \in \mathcal{F}} \chi_f^h.$$

Let the columns of the matrix Φ_Γ be the coefficient vectors of the partition of unity functions; then, the matrix Φ_Γ has only entries 0 and 1. We extend the interface values to the interior using discrete harmonic extensions. The discrete harmonic extension $w := \mathcal{H}_{\Gamma \rightarrow \Omega}(\tau_\Gamma)$ of a finite element function τ_Γ on the interface with respect to the bilinear form $a_\Omega(\cdot, \cdot)$ is given by

$$(3) \quad \begin{aligned} a_{\Omega_l}(w, v) &= 0 \quad \forall v \in V_0^h(\Omega_l), l = 1, \dots, N, \\ w &= \tau_\Gamma \quad \text{on } \Gamma. \end{aligned}$$

Note that a discrete harmonic extension is energy-minimal, i.e., it is

$$a_\Omega(\mathcal{H}_{\Gamma \rightarrow \Omega}(\tau_\Gamma), \mathcal{H}_{\Gamma \rightarrow \Omega}(\tau_\Gamma)) \leq a_\Omega(v, v) \quad \forall v \in \{v \in V^h(\Omega) : v|_\Gamma = \tau_\Gamma\};$$

see, e.g., [53, Sect. 4.4]. In matrix form, the discrete harmonic extension of Φ_Γ can be computed as

$$\Phi = \begin{bmatrix} \Phi_I \\ \Phi_\Gamma \end{bmatrix} = \begin{bmatrix} -K_{II}^{-1} K_{\Gamma I}^T \Phi_\Gamma \\ \Phi_\Gamma \end{bmatrix}.$$

Typically, $K_{II}^{-1} K_{\Gamma I}^T$ is not built explicitly but evaluated from right to left in the application of $K_{II}^{-1} K_{\Gamma I}^T \Phi_\Gamma$. The matrix $K_{II} = \text{diag}_{i=1}^N(K_{II}^{(i)})$ is a block diagonal and contains only the local matrices $K_{II}^{(i)}$ from the nonoverlapping subdomains. Therefore, the factorization of K_{II} can be computed block-by-block and in parallel.

The condition number estimate for the GDSW preconditioner

$$\kappa(M_{\text{GDSW}}^{-1} K) \leq C \left(1 + \frac{H}{\delta}\right) \left(1 + \log\left(\frac{H}{h}\right)\right)^2,$$

cf. [6, 7], holds also for the general case of Ω decomposed into John domains (in two dimensions), and thus, in particular, for unstructured domain decompositions. Note that, in general, the constant C depends on the contrast of the coefficient function A . As a remedy, we will employ the eigenmodes of local generalized eigenvalue problems to compute an adaptive coarse space that is robust, independent of the coefficient function.

5. Adaptive GDSW. In this section, we will introduce the adaptive GDSW (AGDSW) coarse space. Note that we have improved the AGDSW coarse space compared to the variant introduced in [23] for two dimensions. In particular, the

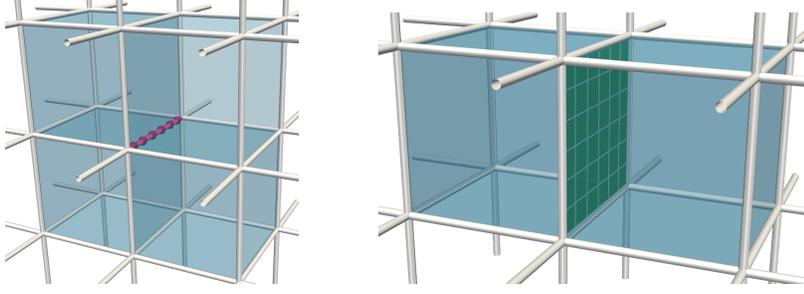


FIG. 2. Graphical representation of $\Omega_e = \bigcup_{k \in n^e} \Omega_k$, the union of all subdomains adjacent to an open edge (left), and $\Omega_f = \Omega_i \cup \Omega_j$, the union of all subdomains adjacent to an open face (right).

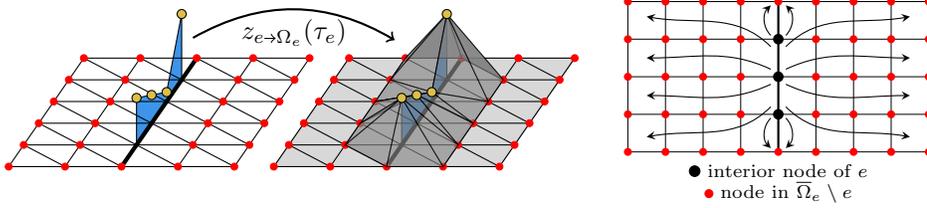


FIG. 3. Graphical representation in two dimensions of the extension by zero of a finite element function defined on an edge $e \in \mathcal{E}$, from the interior degrees of freedom of the edge to the adjacent subdomains (left). Graphical representation in two dimensions of the discrete harmonic extension (4) from the interior degrees of freedom of an edge $e \in \mathcal{E}$ to Ω_e (right).

construction of the eigenvalue problems was simplified. For more details, see [section 7](#). We remark that the proofs for the condition number estimate for the two-dimensional case and for the variant introduced in [23] are analogous to the proof presented here for the 3D case.

5.1. Construction of the AGDSW coarse space. First, we will introduce a generic local generalized eigenvalue problem which is set up for any interface component, i.e., for any edge or face. The coarse basis functions are then constructed as discrete harmonic extensions of corresponding eigenmodes.

Let ξ be an (open) edge e or (open) face f . We denote the set of indices of adjacent subdomains by n^e , n^f , and n^ξ , respectively. Then, we define the set Ω_ξ as the union of all adjacent subdomains; cf. [Figure 2](#). Additionally, we define the following extension-by-zero operator from ξ to a connected set $G \subset \Omega$ with $\xi \subset G$:

$$z_{\xi \rightarrow G} : V^h(\xi) \rightarrow V_0^h(G), \quad v \mapsto z_{\xi \rightarrow G}(v) := \begin{cases} v & \text{in all interior nodes of } \xi, \\ 0 & \text{on all other nodes in } G; \end{cases}$$

see [Figure 3](#) (left) for a graphical representation in 2D. Here,

$$V_0^h(G) := \{v|_G : v \in V^h(\Omega), v = 0 \text{ in } \Omega \setminus G\}.$$

By $\mathcal{H}_{\xi \rightarrow \Omega_\xi}$ we denote the discrete harmonic extension w.r.t. $a_\Omega(\cdot, \cdot)$ from ξ to Ω_ξ . Specifically, let $V_{0,\xi}^h(\Omega_l) := \{w|_{\Omega_l} : w \in V^h(\Omega), w = 0 \text{ on } \xi\}$. Then, for $\tau_\xi \in V^h(\xi)$ the extension $v_\xi := \mathcal{H}_{\xi \rightarrow \Omega_\xi}(\tau_\xi)$ is given by the solution of

$$(4) \quad \begin{aligned} a_{\Omega_l}(v_\xi, v) &= 0 \quad \forall v \in V_{0,\xi}^h(\Omega_l), l \in n^\xi, \\ v_\xi &= \tau_\xi \quad \text{on } \xi, \end{aligned}$$

where n^ξ is the set of indices of all subdomains adjacent to the edge or face ξ . Note that, in contrast to [\(3\)](#), we do not prescribe Dirichlet boundary values on $\Gamma \setminus \xi$.

In particular, the boundary nodes of ξ are part of the Neumann boundary of the discrete harmonic extension $\mathcal{H}_{\xi \rightarrow \Omega_\xi}$; cf. [Figure 3](#) (right). This is different from [\[23\]](#), where finite elements adjacent to the vertices are removed; also see [\[23, Fig. 1\]](#). Our new approach allows to construct the left hand side of the eigenvalue problem from the assembled subdomain stiffness matrix.

Now, we consider the generalized eigenvalue problem on each edge or face ξ : find $\tau_{*,\xi} \in V_0^h(\xi) := \{v|_\xi : v \in V^h(\Omega), v = 0 \text{ on } \partial\xi\}$ such that

$$(5) \quad a_{\Omega_\xi}(\mathcal{H}_{\xi \rightarrow \Omega_\xi}(\tau_{*,\xi}), \mathcal{H}_{\xi \rightarrow \Omega_\xi}(\theta)) = \lambda_{*,\xi} a_{\Omega_\xi}(z_{\xi \rightarrow \Omega_\xi}(\tau_{*,\xi}), z_{\xi \rightarrow \Omega_\xi}(\theta)) \quad \forall \theta \in V_0^h(\xi).$$

Let the eigenvalues be sorted in non-descending order, i.e., $\lambda_{1,\xi} \leq \lambda_{2,\xi} \leq \dots \leq \lambda_{m,\xi}$, and the eigenmodes accordingly, where $m = \dim(V_0^h(\xi))$. Furthermore, let the eigenmodes $\tau_{*,\xi}$ satisfy $a_{\Omega_\xi}(z_{\xi \rightarrow \Omega_\xi}(\tau_{k,\xi}), z_{\xi \rightarrow \Omega_\xi}(\tau_{j,\xi})) = \delta_{kj}$, where δ_{kj} is the Kronecker delta symbol. We select all eigenmodes $\tau_{*,\xi}$ where the eigenvalues are below a certain threshold, i.e., $\lambda_{*,e} \leq \text{tol}_\mathcal{E}$ for edges and $\lambda_{*,f} \leq \text{tol}_\mathcal{F}$ for faces. Then, the coarse basis functions corresponding to ξ are the extensions

$$(6) \quad v_{*,\xi} := \mathcal{H}_{\Gamma \rightarrow \Omega}(z_{\xi \rightarrow \Gamma}(\tau_{*,\xi}))$$

of the selected $\tau_{*,\xi}$.

We define the space of edge based coarse functions as

$$(7) \quad V_\mathcal{E}^{\text{tol}_\mathcal{E}} := \left(\bigoplus_{e \in \mathcal{E}} \text{span} \{v_{k,e} : \lambda_{k,e} \leq \text{tol}_\mathcal{E}\} \right).$$

and the space of face based coarse functions as

$$(8) \quad V_\mathcal{F}^{\text{tol}_\mathcal{F}} := \left(\bigoplus_{f \in \mathcal{F}} \text{span} \{v_{k,f} : \lambda_{k,f} \leq \text{tol}_\mathcal{F}\} \right).$$

In addition to the edge and face basis functions, we use the vertex basis functions

$$\phi_v = \begin{bmatrix} -K_{II}^{-1} K_{\Gamma I}^T \chi_v^h|_\Gamma \\ \chi_v^h|_\Gamma \end{bmatrix}$$

from the GDSW coarse space, and denote the corresponding space by

$$(9) \quad V_\mathcal{V} := \bigoplus_{v \in \mathcal{V}} \text{span} \{\phi_v\};$$

see also [section 4](#). Finally, we obtain the adaptive GDSW coarse space

$$V_{\text{AGDSW}}^{\text{tol}_\mathcal{E}, \text{tol}_\mathcal{F}} = V_\mathcal{V} \oplus V_\mathcal{E}^{\text{tol}_\mathcal{E}} \oplus V_\mathcal{F}^{\text{tol}_\mathcal{F}}.$$

Note that the left hand side of the eigenvalue problem [\(5\)](#) is singular, and its kernel contains the constant functions. Thus, the coarse basis functions corresponding to the eigenvalue 0 are, in fact, the classical coarse GDSW edge and face basis functions. Since $\text{tol}_\mathcal{E}, \text{tol}_\mathcal{F} \geq 0$, these are always included in the adaptive GDSW coarse space.

REMARK 1. For $\text{tol}_\mathcal{E} \geq 0, \text{tol}_\mathcal{F} \geq 0$, we obtain

$$V_{\text{GDSW}} = V_{\text{AGDSW}}^{0,0} \subset V_{\text{AGDSW}}^{\text{tol}_\mathcal{E}, \text{tol}_\mathcal{F}}.$$

REMARK 2. If a Dirichlet boundary condition on $\partial\Omega$ is prescribed only on a subset $\partial\Omega_D \subset \partial\Omega$, in combination with a Neumann boundary condition on $\partial\Omega_N = \partial\Omega \setminus \partial\Omega_D$, the construction of the adaptive GDSW coarse space and the proof of the condition number estimate in [section 6](#) are essentially the same. Finite element nodes that lie on the Neumann boundary are simply treated as interior nodes.

5.2. Properties of the spectral projection. For the coarse interpolation defined in [section 6](#), we consider the projections

$$(10) \quad \Pi_{\mathcal{E}} w := \sum_{e \in \mathcal{E}} \Pi_e w, \quad \Pi_e w := \sum_{\lambda_{k,e} \leq \text{tol}_{\mathcal{E}}} a_{\Omega_e}(z_{e \rightarrow \Omega_e}(w), z_{e \rightarrow \Omega_e}(v_{k,e})) v_{k,e},$$

$$(11) \quad \Pi_{\mathcal{F}} w := \sum_{f \in \mathcal{F}} \Pi_f w, \quad \Pi_f w := \sum_{\lambda_{k,f} \leq \text{tol}_{\mathcal{F}}} a_{\Omega_f}(z_{f \rightarrow \Omega_f}(w), z_{f \rightarrow \Omega_f}(v_{k,f})) v_{k,f}$$

onto the spaces $V_{\mathcal{E}}^{\text{tol}_{\mathcal{E}}}$ and $V_{\mathcal{F}}^{\text{tol}_{\mathcal{F}}}$, respectively. Here, $v_{k,e}$ and $v_{k,f}$ are from [\(6\)](#).

These projections have typical properties, summarized in the following lemma. The lemma may be applied to the projections in [\(10\)](#) and [\(11\)](#); cf. [Lemma 2](#) and [Remark 3](#). The proof uses standard arguments from spectral theory.

LEMMA 1. *Let a symmetric, positive semi-definite bilinear form $d(\cdot, \cdot)$ and a symmetric positive definite bilinear form $c(\cdot, \cdot)$ be given on a finite element space W and consider the eigenvalue problem: find $v \in W$ such that*

$$(12) \quad d(v, w) = \lambda c(v, w) \quad \forall w \in W.$$

Let the corresponding eigenpairs $\{(v_k, \lambda_k)\}_{k=1}^{\dim(W)}$ be chosen such that $c(v_k, v_j) = \delta_{kj}$, where δ_{kj} is the Kronecker delta symbol. Additionally, we assume that the eigenpairs are sorted in non-descending order w.r.t. the eigenvalues. Given $u \in W$, the operator $\Pi u := \sum_{\lambda_k \leq \text{tol}} c(u, v_k) v_k$ defines a projection which is orthogonal with respect to the bilinear form $d(\cdot, \cdot)$ and therefore $|u|_d^2 = |\Pi u|_d^2 + |u - \Pi u|_d^2$. Here, the semi-norm $|u|_d^2$ is defined as $|u|_d^2 := d(u, u)$. In addition, for $\|u\|_c^2 := c(u, u)$, the estimate holds

$$\|u - \Pi u\|_c^2 \leq \frac{1}{\text{tol}} |u - \Pi u|_d^2.$$

LEMMA 2. *Let $\xi \subset \Gamma$ be an open, connected interface component (e.g. an edge or a face) with adjacent subdomains $\Omega_i, i \in n^\xi$. Given a symmetric positive definite bilinear form $c: V^h(\xi) \times V^h(\xi) \rightarrow \mathbb{R}$, assume there exists a constant $C_{\text{inv}, \xi}$,*

$$(13) \quad \text{s.t.} \quad |z_{\xi \rightarrow \Omega_\xi}(v)|_{a, \Omega_\xi}^2 \leq C_{\text{inv}, \xi} \|v\|_c^2 \quad \forall v \in V^h(\xi).$$

Furthermore, let $d: V^h(\xi) \times V^h(\xi) \rightarrow \mathbb{R}$ be a symmetric positive semi-definite bilinear form which satisfies

$$(14) \quad |z_{\xi \rightarrow \Omega_\xi}(v)|_d^2 \leq |v|_{a, \Omega_\xi}^2 \quad \forall v \in V^h(\Omega).$$

Based on the eigenvalue problem [\(12\)](#) in [Lemma 1](#) with $W := V_0^h(\xi)$, we have for $i \in n^\xi$ and $u \in V^h(\Omega)$

$$|z_{\xi \rightarrow \Omega_\xi}(u) - \Pi z_{\xi \rightarrow \Omega_\xi}(u)|_{a, \Omega_i}^2 \leq |z_{\xi \rightarrow \Omega_\xi}(u) - \Pi z_{\xi \rightarrow \Omega_\xi}(u)|_{a, \Omega_\xi}^2 \leq \frac{C_{\text{inv}, \xi}}{\text{tol}_\xi} \sum_{k \in n^\xi} |u|_{a, \Omega_k}^2.$$

Proof of Lemma 2. Using the assumptions and [Lemma 1](#) (third and fourth inequality), we have

$$\begin{aligned} & |z_{\xi \rightarrow \Omega_\xi}(u) - \Pi z_{\xi \rightarrow \Omega_\xi}(u)|_{a, \Omega_i}^2 \leq |z_{\xi \rightarrow \Omega_\xi}(u) - \Pi z_{\xi \rightarrow \Omega_\xi}(u)|_{a, \Omega_\xi}^2 \\ &= |z_{\xi \rightarrow \Omega_\xi}(z_{\xi \rightarrow \Omega_\xi}(u) - \Pi z_{\xi \rightarrow \Omega_\xi}(u))|_{a, \Omega_\xi}^2 \stackrel{(13)}{\leq} C_{\text{inv}, \xi} \|z_{\xi \rightarrow \Omega_\xi}(u) - \Pi z_{\xi \rightarrow \Omega_\xi}(u)\|_c^2 \\ &\stackrel{\text{Lemma 1}}{\leq} \frac{C_{\text{inv}, \xi}}{\text{tol}} |z_{\xi \rightarrow \Omega_\xi}(u) - \Pi z_{\xi \rightarrow \Omega_\xi}(u)|_d^2 \leq \frac{C_{\text{inv}, \xi}}{\text{tol}} |z_{\xi \rightarrow \Omega_\xi}(u)|_d^2 \\ &\stackrel{(14)}{\leq} \frac{C_{\text{inv}, \xi}}{\text{tol}} |u|_{a, \Omega_\xi}^2 = \frac{C_{\text{inv}, \xi}}{\text{tol}} \sum_{k \in n^\xi} |u|_{a, \Omega_k}^2. \end{aligned}$$

□

REMARK 3. For an edge $\xi = e$ or a face $\xi = f$, the bilinear forms are

$$c(\cdot, \cdot) = a_{\Omega_\xi}(z_{\xi \rightarrow \Omega_\xi}(\cdot), z_{\xi \rightarrow \Omega_\xi}(\cdot)) \quad \text{and} \quad d(\cdot, \cdot) = a_{\Omega_\xi}(\mathcal{H}_{\xi \rightarrow \Omega_\xi}(\cdot), \mathcal{H}_{\xi \rightarrow \Omega_\xi}(\cdot));$$

cf. eigenvalue problem (5). Hence, $C_{\text{inv},e} = C_{\text{inv},f} = 1$ and, due to the energy-minimal property of the discrete harmonic extension,

$$\begin{aligned} d(z_{\xi \rightarrow \Omega_\xi}(v), z_{\xi \rightarrow \Omega_\xi}(v)) &= |\mathcal{H}_{\xi \rightarrow \Omega_\xi}(z_{\xi \rightarrow \Omega_\xi}(v))|_{a, \Omega_\xi}^2 = |\mathcal{H}_{\xi \rightarrow \Omega_\xi}(v)|_{a, \Omega_\xi}^2 \\ &\leq |v|_{a, \Omega_\xi}^2 \quad \forall v \in V^h(\Omega). \end{aligned}$$

Thus, for the adaptive GDSW coarse space, the assumptions of Lemma 2 hold.

In subsection 7.1, we describe a variant of adaptive GDSW for which $C_{\text{inv},e} = C_{\text{inv},f}$ corresponds to the constant from an inverse inequality bounding $|\cdot|_{H^1(T)}$ by $\|\cdot\|_{L^2(T)}$ on a finite element T . Subsequently, in subsection 7.2, we describe a variant with a modified left hand side of the generalized eigenvalue problem. Both variants are covered by Lemma 2 and the proof of the existence of a stable decomposition in Theorem 6.

6. Convergence analysis for the overlapping Schwarz method with the adaptive GDSW space. In this section, we will provide a condition number estimate and a proof of this estimate. Following, e.g., [53], we prove the existence of a stable decomposition. Therefore, we have to provide a suitable coarse interpolation I_0 into the coarse space

$$V_0 := V_{\text{AGDSW}}^{\text{tol}_\mathcal{E}, \text{tol}_\mathcal{F}} = V_{\mathcal{V}} \oplus V_{\mathcal{E}}^{\text{tol}_\mathcal{E}} \oplus V_{\mathcal{F}}^{\text{tol}_\mathcal{F}};$$

see (9), (7), and (8) for a definition of $V_{\mathcal{V}}$, $V_{\mathcal{E}}^{\text{tol}_\mathcal{E}}$, and $V_{\mathcal{F}}^{\text{tol}_\mathcal{F}}$.

We construct the coarse interpolant I_0 from a point-wise interpolation

$$I_{\mathcal{V}}u := \sum_{v \in \mathcal{V}} u(v) \phi_v$$

to the space $V_{\mathcal{V}}$ and from the projections $\Pi_{\mathcal{E}}$ and $\Pi_{\mathcal{F}}$ onto the spaces spanned by the edge and face coarse basis functions, respectively; cf. (10) and (11). In particular, we define the coarse component of the stable decomposition as

$$u_0 := I_0u := I_{\mathcal{V}}u + \Pi_{\mathcal{E}}u + \Pi_{\mathcal{F}}u.$$

The projection operators $I_{\mathcal{V}}$, $\Pi_{\mathcal{E}}$, and $\Pi_{\mathcal{F}}$ satisfy the following assumption:

ASSUMPTION 1. As in Lemma 2, let ξ be an open and connected interface component and Π the corresponding projection operator. Then,

$$z_{\xi \rightarrow \Omega_\xi}(\Pi_*v) = \Pi_*z_{\xi \rightarrow \Omega_\xi}(v) = 0 \quad \forall v \in V^h(\Omega)$$

for any other projection operator $\Pi_* \neq \Pi$.

This assumption is satisfied for the projection operators of the AGDSW coarse space interpolation. In particular, we have

$$\begin{aligned} I_{\mathcal{V}}z_{e \rightarrow \Omega_e}(v) &= z_{e \rightarrow \Omega_e}(I_{\mathcal{V}}v) = \Pi_{\mathcal{F}}z_{e \rightarrow \Omega_e}(v) = z_{e \rightarrow \Omega_e}(\Pi_{\mathcal{F}}v) = 0 \quad \forall v \in V^h(\Omega) \quad \text{and} \\ I_{\mathcal{V}}z_{f \rightarrow \Omega_f}(v) &= z_{f \rightarrow \Omega_f}(I_{\mathcal{V}}v) = \Pi_{\mathcal{E}}z_{f \rightarrow \Omega_f}(v) = z_{f \rightarrow \Omega_f}(\Pi_{\mathcal{E}}v) = 0 \quad \forall v \in V^h(\Omega), \end{aligned}$$

which follows from the definition of $z_{e \rightarrow \Omega_e}(\cdot)$ and $z_{f \rightarrow \Omega_f}(\cdot)$, since vertex basis functions vanish on edges and faces, edge basis functions vanish on vertices and faces, and face basis functions vanish on vertices and edges.

To prove the existence of a stable decomposition, we first prove the following lemma. It states estimates for the edge and face functions that arise during the proof; cf. Theorem 6.

LEMMA 3. Let the assumptions of Lemma 2 and Assumption 1 be satisfied, then

$$|z_{\xi \rightarrow \Omega_{\xi}}(u - u_0)|_{a, \Omega_i}^2 \leq |z_{\xi \rightarrow \Omega_{\xi}}(u - u_0)|_{a, \Omega_{\xi}}^2 \leq \frac{C_{inv, \xi}}{tol_{\xi}} \sum_{k \in n^{\xi}} |u|_{a, \Omega_k}^2$$

for an edge $\xi = e \in \mathcal{E}$ or a face $\xi = f \in \mathcal{F}$.

Proof. Due to Assumption 1 we have on an edge e

$$\begin{aligned} z_{e \rightarrow \Omega_e}(u - u_0) &= z_{e \rightarrow \Omega_e}(u) - z_{e \rightarrow \Omega_e}(I_{\mathcal{V}}u + \Pi_{\mathcal{E}}u + \Pi_{\mathcal{F}}u) \\ &\stackrel{A_{sm}.1}{=} z_{e \rightarrow \Omega_e}(u) - z_{e \rightarrow \Omega_e}(\Pi_{\mathcal{E}}u) = z_{e \rightarrow \Omega_e}(u) - z_{e \rightarrow \Omega_e}(\Pi_e u) \\ &= z_{e \rightarrow \Omega_e}(u) - \Pi_e z_{e \rightarrow \Omega_e}(u) \end{aligned}$$

and, analogously, $z_{f \rightarrow \Omega_f}(u - u_0) = z_{f \rightarrow \Omega_f}(u) - \Pi_f z_{f \rightarrow \Omega_f}(u)$ on a face f . Therefore, using Lemma 2 we obtain

$$\begin{aligned} |z_{\xi \rightarrow \Omega_{\xi}}(u - u_0)|_{a, \Omega_i}^2 &\leq |z_{\xi \rightarrow \Omega_{\xi}}(u - u_0)|_{a, \Omega_{\xi}}^2 = |z_{\xi \rightarrow \Omega_{\xi}}(u) - \Pi_{\xi} z_{\xi \rightarrow \Omega_{\xi}}(u)|_{a, \Omega_{\xi}}^2 \\ &\stackrel{Lemma\ 2}{\leq} \frac{C_{inv, \xi}}{tol_{\xi}} \sum_{k \in n^{\xi}} |u|_{a, \Omega_k}^2. \end{aligned}$$

Next, we derive an estimate for the energy of the coarse component on a subdomain.

LEMMA 4. Under the assumptions of Lemma 3, for $i \in \{1, \dots, N\}$, we have

$$|u_0|_{a, \Omega_i}^2 \leq 2|u|_{a, \Omega_i}^2 + \frac{4N^e C_{inv, e}}{tol_{\mathcal{E}}} \sum_{e \subset \partial \Omega_i} \sum_{k \in n^e} |u|_{a, \Omega_k}^2 + \frac{4N^f C_{inv, f}}{tol_{\mathcal{F}}} \sum_{f \subset \partial \Omega_i} \sum_{k \in n^f} |u|_{a, \Omega_k}^2,$$

where N^e and N^f denote the maximum number of edges and faces, respectively, a subdomain can have.

Proof. We can use the fact that u_0 is discrete harmonic on each subdomain Ω_i and consider the contributions on the interface components separately. Since $u - u_0 = 0$ in the vertices, we obtain

$$\begin{aligned} |u_0|_{a, \Omega_i}^2 &\leq 2|\mathcal{H}_{\partial \Omega_i \rightarrow \Omega_i}(u)|_{a, \Omega_i}^2 + 2|\mathcal{H}_{\partial \Omega_i \rightarrow \Omega_i}(u - u_0)|_{a, \Omega_i}^2 \\ &= 2(|\mathcal{H}_{\partial \Omega_i \rightarrow \Omega_i}(u)|_{a, \Omega_i}^2 \\ &\quad + \left| \sum_{e \subset \partial \Omega_i} \mathcal{H}_{\partial \Omega_i \rightarrow \Omega_i}(z_{e \rightarrow \partial \Omega_i}(u - u_0)) + \sum_{f \subset \partial \Omega_i} \mathcal{H}_{\partial \Omega_i \rightarrow \Omega_i}(z_{f \rightarrow \partial \Omega_i}(u - u_0)) \right|_{a, \Omega_i}^2) \end{aligned}$$

Using a Cauchy-Schwarz inequality and the energy-minimality of $\mathcal{H}_{\partial \Omega_i \rightarrow \Omega_i}(\cdot)$ gives

$$\begin{aligned} |u_0|_{a, \Omega_i}^2 &\leq 2|u|_{a, \Omega_i}^2 + 4N^e \sum_{e \subset \partial \Omega_i} |z_{e \rightarrow \Omega_e}(u - u_0)|_{a, \Omega_i}^2 \\ &\quad + 4N^f \sum_{f \subset \partial \Omega_i} |z_{f \rightarrow \Omega_f}(u - u_0)|_{a, \Omega_i}^2, \end{aligned}$$

where N^e and N^f denote the maximum number of edges and faces, respectively, a subdomain can have. Finally, using Lemma 3, we obtain

$$|u_0|_{a, \Omega_i}^2 \leq 2|u|_{a, \Omega_i}^2 + \frac{4N^e C_{inv, e}}{tol_{\mathcal{E}}} \sum_{e \subset \partial \Omega_i} \sum_{k \in n^e} |u|_{a, \Omega_k}^2 + \frac{4N^f C_{inv, f}}{tol_{\mathcal{F}}} \sum_{f \subset \partial \Omega_i} \sum_{k \in n^f} |u|_{a, \Omega_k}^2. \quad \square$$

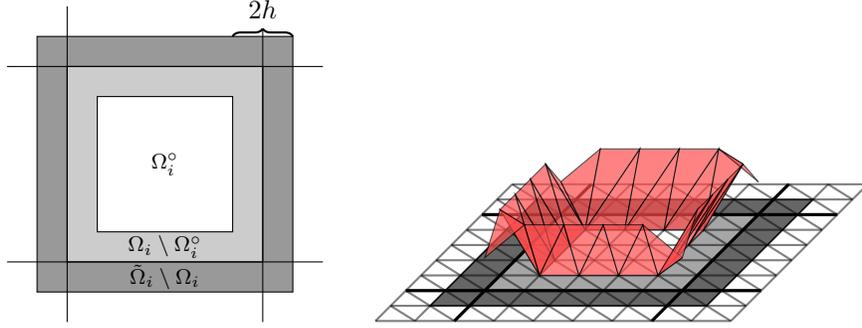


FIG. 4. In the proof of [Theorem 6](#), we consider a partition of unity corresponding to an overlapping decomposition $\{\tilde{\Omega}_i\}_{i=1}^N$ with overlap h . The corresponding regions Ω_i° (white), $G_1 := \Omega_i \setminus \Omega_i^\circ$ (light gray), and $G_2 := \tilde{\Omega}_i \setminus \Omega_i$ (dark gray) are depicted in two dimensions (**left**). A finite element function ψ on G that is constant on the edges and vanishes on $\partial\Omega_i^\circ$ (**right**); cf. [Lemma 5](#). Left figure from [[24](#), Fig. 6.1].

In [Theorem 6](#), we prove the existence of a stable decomposition by introducing an overlapping decomposition $\{\tilde{\Omega}_i\}_{i=1}^N$ with overlap h corresponding to the nonoverlapping decomposition $\{\Omega_i\}_{i=1}^N$; cf. [Figure 4](#). Note that, in general, this decomposition differs from $\{\Omega'_i\}_{i=1}^N$, the one used in the first level of the preconditioner with overlap $\delta \geq h$. The decomposition $\{\tilde{\Omega}_i\}_{i=1}^N$ is only used in the proof and does not place any restriction on δ . However, it does remove the dependence of the condition number estimate in [Corollary 7](#) on the size of the overlap δ .

In the proof of the stable decomposition, we need to estimate the energy of local components $I^h(\theta_i(u - u_0))$, given a partition of unity $\{\theta_i\}_{i=1}^N$ defined on the overlapping decomposition $\{\tilde{\Omega}_i\}_{i=1}^N$. The following lemma is used in [Theorem 6](#) to estimate the energy of the local components on the overlap.

LEMMA 5. *Let the assumptions of [Lemma 3](#) be satisfied. For any subdomain Ω_i , $i \in \{1, \dots, N\}$, let $G \subseteq \tilde{\Omega}_i \setminus \Omega_i^\circ$, where Ω_i° denotes the non-overlapping subset of Ω_i ; cf. [Figure 4](#). We consider a finite element function $\psi \in V^h(G)$ that can have arbitrary values on $\partial\Omega_i$ but vanishes on $\partial(\tilde{\Omega}_i \setminus \Omega_i^\circ) \cap \bar{G}$. Moreover, we assume that $0 \leq \psi \leq 1$, and that $\psi|_e$ and $\psi|_f$ are constant on $e \in \mathcal{E}$ and $f \in \mathcal{F}$, respectively. Then,*

$$\left| I^h(\psi \cdot (u - u_0)) \right|_{a,G}^2 \leq \frac{2N^e C_{\text{inv},e}}{\text{tol}_{\mathcal{E}}} \sum_{e \in \partial\Omega_i} \sum_{k \in n^e} |u|_{a,\Omega_k}^2 + \frac{2N^f C_{\text{inv},f}}{\text{tol}_{\mathcal{F}}} \sum_{f \subset \partial\Omega_i} \sum_{k \in n^f} |u|_{a,\Omega_k}^2,$$

where N^e and N^f correspond to the maximum number of edges and faces, respectively, a subdomain can have.

Note that in the proof of the stable decomposition in [Theorem 6](#), we will make use of [Lemma 5](#) with the sets $G = G_1$ and $G = G_2$; cf. [Figure 4](#).

Proof. We observe that $z_{e \rightarrow \Omega}(\cdot)$ and $z_{f \rightarrow \Omega}(\cdot)$ are identity operators on e and f , respectively, and that $u - u_0$ vanishes in the vertices. Then, since $G_1 = \Omega_i \setminus \Omega_i^\circ$ and $G_2 = \tilde{\Omega}_i \setminus \Omega_i$ have width $1h$ and since ψ , $z_{e \rightarrow \Omega_e}$, and $z_{f \rightarrow \Omega_f}$ all vanish on $\partial(\tilde{\Omega}_i \setminus \Omega_i^\circ) \cap \bar{G}$, we have with $\psi_0 := \psi \cdot (u - u_0)$

$$\begin{aligned} |I^h(\psi_0)|_{a,G}^2 &= \left| \sum_{e \subset \partial\Omega_i} z_{e \rightarrow \Omega}(\psi_0) + \sum_{f \subset \partial\Omega_i} z_{f \rightarrow \Omega}(\psi_0) \right|_{a,G}^2 \\ &\leq 2 \left| \sum_{e \subset \partial\Omega_i} z_{e \rightarrow \Omega}(\psi_0) \right|_{a,G}^2 + 2 \left| \sum_{f \subset \partial\Omega_i} z_{f \rightarrow \Omega}(\psi_0) \right|_{a,G}^2. \end{aligned}$$

Then, a Cauchy-Schwarz inequality, the fact that $\psi|_e$ is constant and that $0 \leq \psi \leq 1$, gives

$$\begin{aligned} \left| \sum_{e \in \partial\Omega_i} z_{e \rightarrow \Omega} \left(\psi(u - u_0) \right) \right|_{a,G}^2 &\leq N^e \sum_{e \in \partial\Omega_i} \left| z_{e \rightarrow \Omega} \left(\psi(u - u_0) \right) \right|_{a,G}^2 \\ &\leq N^e \sum_{e \in \partial\Omega_i} (\psi|_e)^2 |z_{e \rightarrow \Omega}(u - u_0)|_{a,G}^2 \leq N^e \sum_{e \in \partial\Omega_i} |z_{e \rightarrow \Omega}(u - u_0)|_{a,G}^2. \end{aligned}$$

Finally, using [Lemma 3](#), we obtain

$$\sum_{e \in \partial\Omega_i} |z_{e \rightarrow \Omega}(u - u_0)|_{a,G}^2 \leq \sum_{e \in \partial\Omega_i} |z_{e \rightarrow \Omega}(u - u_0)|_{a,\Omega_e}^2 \leq \frac{C_{\text{inv},e}}{\text{tol}_{\mathcal{E}}} \sum_{e \in \partial\Omega_i} \sum_{k \in n^e} |u|_{a,\Omega_k}^2.$$

Completely analogously, using [Lemma 3](#), we have

$$\left| \sum_{f \in \partial\Omega_i} z_{f \rightarrow \Omega} \left(\psi(u - u_0) \right) \right|_{a,G}^2 \leq \frac{N^f C_{\text{inv},f}}{\text{tol}_{\mathcal{F}}} \sum_{f \in \partial\Omega_i} \sum_{k \in n^f} |u|_{a,\Omega_k}^2.$$

Therefore,

$$\begin{aligned} |I^h(\psi_0)|_{a,G}^2 &\leq 2 \left| \sum_{e \in \partial\Omega_i} z_{e \rightarrow \Omega}(\psi_0) \right|_{a,G}^2 + 2 \left| \sum_{f \in \partial\Omega_i} z_{f \rightarrow \Omega}(\psi_0) \right|_{a,G}^2 \\ &\leq \frac{2N^e C_{\text{inv},e}}{\text{tol}_{\mathcal{E}}} \sum_{e \in \partial\Omega_i} \sum_{k \in n^e} |u|_{a,\Omega_k}^2 + \frac{2N^f C_{\text{inv},f}}{\text{tol}_{\mathcal{F}}} \sum_{f \in \partial\Omega_i} \sum_{k \in n^f} |u|_{a,\Omega_k}^2. \quad \square \end{aligned}$$

Now, we are able to prove the existence of a stable decomposition.

THEOREM 6 (Stable Decomposition). *Under the assumptions of [Lemma 3](#), for each $u \in V^h(\Omega)$, there exists a decomposition $u = \sum_{i=0}^N R_i^T u_i$, $u_i \in V_i = V^h(\Omega'_i)$, where $\Omega'_0 := \Omega$, such that*

$$\sum_{i=0}^N |u_i|_{a,\Omega'_i}^2 \leq C_0^2 |u|_{a,\Omega}^2,$$

where $C_0^2 = \left(20 + \frac{34(N^e)^2 n_{\text{max}}^e C_{\text{inv},e}}{\text{tol}_{\mathcal{E}}} + \frac{68(N^f)^2 C_{\text{inv},f}}{\text{tol}_{\mathcal{F}}} \right)$ and N^e and N^f correspond to the maximum number of edges and faces, respectively, a subdomain can have, and n_{max}^e corresponds to the maximum number of adjacent subdomains an edge can have.

Proof. On the overlapping decomposition $\{\tilde{\Omega}_i\}_{i=1}^N$ of width $1h$, we consider the local components $u_i := I^h(\theta_i(u - u_0))$ with the partition of unity $\{\theta_i\}_{i=1}^N$, $\theta_i \in V^h(\Omega)$, where

$$\theta_i(x^h) := \begin{cases} \frac{1}{|n^e|} & \text{on edges } e \in \mathcal{E}, \\ \frac{1}{|n^f|} & \text{on faces } f \in \mathcal{F}, \\ \frac{1}{|n^v|} & \text{on vertices } v \in \mathcal{V}, \\ 1 & \text{in } \Omega_i^\circ, \\ 0 & \text{elsewhere,} \end{cases}$$

where x^h is a finite element node and where Ω_i° denotes the non-overlapping subset of $\tilde{\Omega}_i$; cf. [Figure 4](#). Note that, since $\tilde{\Omega}_i \subset \Omega'_i$, we still have $u_i \in V_i$. We consider the partition

$$\tilde{\tilde{\Omega}}_i = (\tilde{\Omega}_i \setminus \Omega_i) \cup (\Omega_i \setminus \Omega_i^\circ) \cup \Omega_i^\circ.$$

Therefore, we have

$$(15) \quad |u_i|_{a,\tilde{\Omega}_i}^2 = |u_i|_{a,\tilde{\Omega}_i \setminus \Omega_i}^2 + |u_i|_{a,\Omega_i \setminus \Omega_i^\circ}^2 + |u_i|_{a,\Omega_i^\circ}^2.$$

To proceed, we use the estimate for $|u_0|_{a,\Omega_i}^2$ from [Lemma 4](#). Let

$$Z_{\mathcal{E}} := \frac{4N^e C_{\text{inv},e}}{\text{tol}_{\mathcal{E}}} \sum_{e \subset \partial \Omega_i} \sum_{k \in n^e} |u|_{a,\Omega_k}^2, \quad Z_{\mathcal{F}} := \frac{4N^f C_{\text{inv},f}}{\text{tol}_{\mathcal{F}}} \sum_{f \subset \partial \Omega_i} \sum_{k \in n^f} |u|_{a,\Omega_k}^2,$$

then we have for the last additive term in [\(15\)](#)

$$(16) \quad \begin{aligned} |u_i|_{a,\Omega_i^\circ}^2 &= |I^h(\theta_i(u - u_0))|_{a,\Omega_i^\circ}^2 = |u - u_0|_{a,\Omega_i^\circ}^2 \leq |u - u_0|_{a,\Omega_i}^2 \\ &\leq 2|u|_{a,\Omega_i}^2 + 2|u_0|_{a,\Omega_i}^2 \stackrel{\text{Lemma 4}}{\leq} 2|u|_{a,\Omega_i}^2 + 2\left(2|u|_{a,\Omega_i}^2 + Z_{\mathcal{E}} + Z_{\mathcal{F}}\right) \\ &= 6|u|_{a,\Omega_i}^2 + 2Z_{\mathcal{E}} + 2Z_{\mathcal{F}}. \end{aligned}$$

Furthermore, we have for the second additive term in [\(15\)](#)

$$(17) \quad \begin{aligned} |u_i|_{a,\Omega_i \setminus \Omega_i^\circ}^2 &\leq 2|u_i - (u - u_0)|_{a,\Omega_i \setminus \Omega_i^\circ}^2 + 2|u - u_0|_{a,\Omega_i \setminus \Omega_i^\circ}^2 \\ &\leq 2|I^h((1 - \theta_i)(u - u_0))|_{a,\Omega_i \setminus \Omega_i^\circ}^2 + 2|u - u_0|_{a,\Omega_i}^2. \end{aligned}$$

We observe that, on an edge $e \in \mathcal{E}$ or a face $f \in \mathcal{F}$, the restrictions of θ_i are constant according to its definition:

$$\theta_i|_e = \frac{1}{|n^e|} \leq \frac{1}{2}, \quad \theta_i|_f = \frac{1}{|n^f|} = \frac{1}{2}.$$

Therefore, setting $\psi := 1 - \theta_i$ and $G := G_1 = \Omega_i \setminus \Omega_i^\circ$, we can use [Lemma 5](#) to bound the first additive term of equation [\(17\)](#). Note that we cannot set $\psi = \theta_i$ to derive an estimate for $|u_i|_{a,\Omega_i \setminus \Omega_i^\circ}^2$ directly, since [Lemma 5](#) requires $\psi = 0$ on the boundary of Ω_i° . Using [Lemma 4](#) and equation [\(16\)](#), we obtain for equation [\(17\)](#)

$$(18) \quad \begin{aligned} |u_i|_{a,\Omega_i \setminus \Omega_i^\circ}^2 &\leq 2|I^h((1 - \theta_i)(u - u_0))|_{a,\Omega_i \setminus \Omega_i^\circ}^2 + 2|u - u_0|_{a,\Omega_i}^2 \\ &\leq 2\left(0.5Z_{\mathcal{E}} + 0.5Z_{\mathcal{F}}\right) + 2\left(6|u|_{a,\Omega_i}^2 + 2Z_{\mathcal{E}} + 2Z_{\mathcal{F}}\right) \\ &= 12|u|_{a,\Omega_i}^2 + 5Z_{\mathcal{E}} + 5Z_{\mathcal{F}}. \end{aligned}$$

Now, setting $\psi := \theta_i$ on $G := G_2 = \tilde{\Omega}_i \setminus \Omega_i$ and using [Lemma 5](#), we have

$$(19) \quad |u_i|_{a,\tilde{\Omega}_i \setminus \Omega_i}^2 = \left|I^h(\theta_i(u - u_0))\right|_{a,\tilde{\Omega}_i \setminus \Omega_i}^2 \leq 0.5Z_{\mathcal{E}} + 0.5Z_{\mathcal{F}}.$$

Summing the edge and face contributions $Z_{\mathcal{E}}$ and $Z_{\mathcal{F}}$ over all subdomains, we obtain

$$(20) \quad \begin{aligned} \sum_{i=1}^N (Z_{\mathcal{E}} + Z_{\mathcal{F}}) &= \sum_{i=1}^N 4 \left(\frac{N^e C_{\text{inv},e}}{\text{tol}_{\mathcal{E}}} \sum_{e \subset \partial \Omega_i} \sum_{k \in n^e} |u|_{a,\Omega_k}^2 + \frac{N^f C_{\text{inv},f}}{\text{tol}_{\mathcal{F}}} \sum_{f \subset \partial \Omega_i} \sum_{k \in n^f} |u|_{a,\Omega_k}^2 \right) \\ &\leq \frac{4(N^e)^2 n_{\text{max}}^e C_{\text{inv},e}}{\text{tol}_{\mathcal{E}}} |u|_{a,\Omega}^2 + \frac{4(N^f)^2 2C_{\text{inv},f}}{\text{tol}_{\mathcal{F}}} |u|_{a,\Omega}^2, \end{aligned}$$

where n_{max}^e corresponds to the maximum number of adjacent subdomains of an edge. Finally, using

$$|u_0|_{a,\Omega}^2 = \sum_{i=1}^N |u_0|_{a,\Omega_i}^2,$$

we obtain with [Lemma 4](#) and equations [\(16\)](#), [\(18\)](#), [\(19\)](#), and [\(20\)](#)

$$\begin{aligned} \sum_{i=0}^N |u_i|_{a,\Omega}^2 &= \sum_{i=1}^N \left(|u_0|_{a,\Omega_i}^2 + |u_i|_{a,\bar{\Omega}_i \setminus \Omega_i}^2 + |u_i|_{a,\Omega_i \setminus \Omega_i^c}^2 + |u_i|_{a,\Omega_i^c}^2 \right) \\ &\leq \sum_{i=1}^N \left(20 |u|_{a,\Omega_i}^2 + 8.5 Z_{\mathcal{E}} + 8.5 Z_{\mathcal{F}} \right) \\ &\leq \left(20 + \frac{34(N^e)^2 n_{\max}^e C_{\text{inv},e}}{\text{tol}_{\mathcal{E}}} + \frac{68(N^f)^2 C_{\text{inv},f}}{\text{tol}_{\mathcal{F}}} \right) |u|_{a,\Omega}^2. \quad \square \end{aligned}$$

From [Theorem 6](#), we directly obtain a condition number estimate for the pre-conditioned system.

COROLLARY 7. *The condition number of the AGDSW two level Schwarz operator in three dimensions is bounded by*

$$\kappa(M_{\text{AGDSW}}^{-1}K) \leq \left(20 + \frac{34(N^e)^2 n_{\max}^e}{\text{tol}_{\mathcal{E}}} + \frac{68(N^f)^2}{\text{tol}_{\mathcal{F}}} \right) (\hat{N}_c + 1).$$

The constant \hat{N}_c is an upper bound for the number of overlapping subdomains each point $x \in \Omega$ can belong to. All constants are independent of H , h , and the contrast of the coefficient function A .

Proof. Since we use exact local solvers, we directly obtain

$$\kappa(M_{\text{AGDSW}}^{-1}K) \leq C_0^2 (\hat{N}_c + 1),$$

where C_0^2 is the constant of the stable decomposition; cf. [\[53, Lemma 3.11\]](#) and the follow-up discussion and the proof of [\[13, Theorem 4.1\]](#). We obtain the final estimate using [Theorem 6](#) and $C_{\text{inv},e} = C_{\text{inv},f} = 1$; cf. [Remark 3](#). \square

REMARK 4. *The proof for the two-dimensional case can be performed analogously to the three-dimensional case. In particular, the edges in two dimensions can be handled in the same way as the faces in three dimensions. For the AGDSW two level Schwarz operator, we obtain the condition number bound*

$$\kappa(M_{\text{AGDSW}}^{-1}K) \leq \left(20 + \frac{68(N^e)^2}{\text{tol}_{\mathcal{E}}} \right) (\hat{N}_c + 1).$$

7. Variants of adaptive GDSW. There are several modifications that can be applied to the AGDSW coarse space. First, the right hand side of the eigenvalue problem can be replaced by a bilinear form that corresponds to a scaled L^2 -inner product or a scaled mass matrix, respectively; cf. [subsection 7.1](#). To the best of our knowledge, this modification does not lead to an advantage, since the mass matrix has to be additionally assembled whereas the stiffness matrices in an implementation of the right hand side of the eigenvalue problem [\(5\)](#) can be extracted directly from the fully assembled global stiffness matrix K . However, it shows the connection of the AGDSW coarse space to other related coarse spaces, e.g., the OS-ACMS, SHEM, wirebasket, and vertex-based coarse spaces; see [\[24, 19, 16\]](#). Second, a parallel implementation of the left hand side of the eigenvalue problem [\(5\)](#) is facilitated and the computational cost is reduced by using a sum of local, decoupled parts that can then be computed independently; cf. [subsection 7.2](#). A third modification can be used to further decrease the work for the computation of the left hand side of the eigenvalue problem. Here, we consider discrete harmonic extensions onto slabs of finite elements instead of the union of all subdomains which are adjacent to the corresponding edge or face; cf. [subsection 7.3](#).

REMARK 5. *The standard AGDSW algorithm and the mentioned modifications above can also be used in two dimensions, see [23], in which the edge basis functions were constructed slightly differently. The construction presented in subsection 5.1 significantly simplifies the setup of the generalized eigenvalue problems, reduces the computational cost, and can decrease the coarse space dimension.*

7.1. Mass matrix. As in other adaptive coarse spaces, where the generalized eigenvalue problem is used to replace a Poincaré type inequality, cf., e.g. [17, 14, 12], we can use a scaled mass matrix on the right hand side of the eigenvalue problems (5) as well. Let ξ be an edge $e \in \mathcal{E}$ or a face $f \in \mathcal{F}$, then the scaled mass matrix corresponding to the edge e or face f arises from the discretization of the scaled L^2 -inner product

$$b_\xi(u, v) := \frac{1}{h^2} \left(A z_{\xi \rightarrow \Omega_\xi}(u), z_{\xi \rightarrow \Omega_\xi}(v) \right)_{L^2(\Omega_\xi)}.$$

The corresponding norm is defined as

$$\|v\|_{b, \xi}^2 := b_\xi(v, v).$$

Therefore, we obtain for the generalized eigenvalue problem: find $\tau_{*, \xi} \in V_0^h(\xi)$, s.t.

$$a_{\Omega_\xi}(\mathcal{H}_{\xi \rightarrow \Omega_\xi}(\tau_{*, \xi}), \mathcal{H}_{\xi \rightarrow \Omega_\xi}(\theta)) = \lambda_{*, \xi} b_\xi(\tau_{*, \xi}, \theta) \quad \forall \theta \in V_0^h(\xi).$$

We denote the resulting coarse space by $V_{\text{AGDSW-M}}$. For $v \in V^h(\xi)$, we have

$$|z_{\xi \rightarrow \Omega_\xi}(v)|_{a, \Omega_\xi}^2 = \int_{\Omega_\xi} A (\nabla z_{\xi \rightarrow \Omega_\xi}(v))^2 dx \leq \frac{C_{\text{inv}}}{h^2} \int_{\Omega_\xi} A z_{\xi \rightarrow \Omega_\xi}(v)^2 dx = b_\xi(v, v),$$

since A is constant on each fine element $T \in \tau_h(\Omega)$. The constant $C_{\text{inv}} > 0$ arises from the use of an inverse equality on the elements. It is independent of H , h , and the contrast of the coefficient function.

REMARK 6. *The constant C_{inv} depends only on the shape parameter of the triangulation and the polynomial degree of the shape functions; see, e.g., [54, Section 3.6], where also a concrete upper bound for C_{inv} is given.*

We obtain a condition number bound analogously to Corollary 7 by setting $c(\cdot, \cdot) := b_\xi(\cdot, \cdot)$ in Lemma 2.

COROLLARY 8. *The condition number of the AGDSW-M two level Schwarz operator in three dimensions is bounded by*

$$\kappa(M_{\text{AGDSW-M}}^{-1}K) \leq \left(20 + 34C_{\text{inv}} \left(\frac{(N^e)^2 n_{\text{max}}^e}{\text{tol}_\mathcal{E}} + \frac{2(N^f)^2}{\text{tol}_\mathcal{F}} \right) \right) (\hat{N}_c + 1).$$

All constants are independent of H , h , and the contrast of the coefficient function.

REMARK 7. *If the mesh regularity or uniformity is low, better numerical results may be achieved by scaling element-wise with the radius of the largest insphere, i.e., let r_s be a function that is constant on each finite element $T \in \tau_h(\Omega)$, on which it assumes the radius of the largest insphere of T . Then, we define on an edge $\xi = e \in \mathcal{E}$ or a face $\xi = f \in \mathcal{F}$*

$$b_\xi(u, v) := \left(\frac{A}{r_s^2} z_{\xi \rightarrow \Omega_\xi}(u), z_{\xi \rightarrow \Omega_\xi}(v) \right)_{L^2(\Omega_\xi)}.$$

7.2. Local Neumann problems. In a parallel implementation of the generalized face eigenvalue problem (5) we can utilize the fact that the discrete harmonic extension is only weakly coupled via the boundary nodes of the face. Thus, instead of computing the (coupled) extension simultaneously to both subdomains adjacent to the face, we can compute extensions independently to each adjacent subdomain without losing much information. The same holds in two dimensions for the edge eigenvalue problems. Similarly, in three dimensions for edges, we can compute the discrete harmonic extensions independently to the adjacent subdomains. However, the stronger the coupling between the subdomains, the more information is lost, which can result in an increased coarse space dimension.

Let either $\xi = e \in \mathcal{E}$ or $\xi = f \in \mathcal{F}$, then

$$a_{\Omega_\xi}(\mathcal{H}_{\xi \rightarrow \Omega_\xi}(\tau_{*,\xi}), \mathcal{H}_{\xi \rightarrow \Omega_\xi}(\theta)) \neq \sum_{k \in n^\xi} a_{\Omega_k}(\mathcal{H}_{\xi \rightarrow \Omega_k}(\tau_{*,\xi}), \mathcal{H}_{\xi \rightarrow \Omega_k}(\theta)).$$

Nevertheless, we can replace the left hand side of the eigenvalue problem by the sum of the local contributions and obtain the eigenvalue problems: find $\tau_{*,\xi} \in V_0^h(\xi)$ s.t.

$$\sum_{k \in n^\xi} a_{\Omega_k}(\mathcal{H}_{\xi \rightarrow \Omega_k}(\tau_{*,\xi}), \mathcal{H}_{\xi \rightarrow \Omega_k}(\theta)) = \lambda_{*,\xi} a_{\Omega_\xi}(z_{\xi \rightarrow \Omega_\xi}(\tau_{*,\xi}), z_{\xi \rightarrow \Omega_\xi}(\theta)) \quad \forall \theta \in V_0^h(\xi).$$

We denote the resulting coarse space by $V_{\text{AGDSW-S}}$. Using these modified eigenvalue problems yields the same condition number estimate as in [Corollary 7](#).

COROLLARY 9. *The condition number of the AGDSW-S two level Schwarz operator in three dimensions is bounded by*

$$\kappa(M_{\text{AGDSW-S}}^{-1}K) \leq \left(20 + \frac{34(N^e)^2 n_{\max}^e}{\text{tol}_\mathcal{E}} + \frac{68(N^f)^2}{\text{tol}_\mathcal{F}}\right) (\hat{N}_c + 1).$$

All constants are independent of H , h , and the contrast of the coefficient function A .

Proof. We only have to show that the assumptions of [Lemma 2](#) are satisfied. Then, the proof is exactly the same as for [Corollary 7](#).

The bilinear form $d(\cdot, \cdot) := \sum_{k \in n^\xi} a_{\Omega_k}(\mathcal{H}_{\xi \rightarrow \Omega_k}(\cdot), \mathcal{H}_{\xi \rightarrow \Omega_k}(\cdot))$ is symmetric and positiv semi-definite and satisfies

$$d(v, v) = \sum_{k \in n^\xi} |\mathcal{H}_{\xi \rightarrow \Omega_k}(v)|_{a, \Omega_k}^2 \leq \sum_{k \in n^\xi} |v|_{a, \Omega_k}^2 = |v|_{a, \Omega_\xi}^2 \quad \forall v \in V^h(\xi). \quad \square$$

As we are going to observe, this variant of AGDSW can lead to a slightly larger coarse space. However, the implementation and computation of the eigenvalue problems is simplified.

7.3. Economic version using slabs. In order to reduce the computational cost of the computation of the eigenvalue problems, the size of the sets Ω_e and Ω_f can be reduced. In particular, we propose a variant where slabs of width l elements around the edges or faces are used instead of complete subdomains; cf. [Figure 5](#) for the graphical representation of the slabs. We denote these slabs by Ω_e^l and Ω_f^l . The idea of computing the Schur complement only on slabs of minimal width was initially proposed in [\[10\]](#). It was then applied to eigenvalue problems and more general slabs in [\[38\]](#). Finally, in [\[24\]](#), a multiscale coarse space based on the ACMS space was introduced for which an economic variant on slabs was proposed.

The modified eigenvalue problem reads for an edge $\xi = e \in \mathcal{E}$ or a face $\xi = f \in \mathcal{F}$: find $\tau_{*,\xi} \in V_0^h(\xi)$ such that

$$a_{\Omega_\xi^l}(\mathcal{H}_{\xi \rightarrow \Omega_\xi^l}(\tau_{*,\xi}), \mathcal{H}_{\xi \rightarrow \Omega_\xi^l}(\theta)) = \lambda_{*,\xi} a_{\Omega_\xi}(z_{\xi \rightarrow \Omega_\xi}(\tau_{*,\xi}), z_{\xi \rightarrow \Omega_\xi}(\theta)) \quad \forall \theta \in V_0^h(\xi).$$

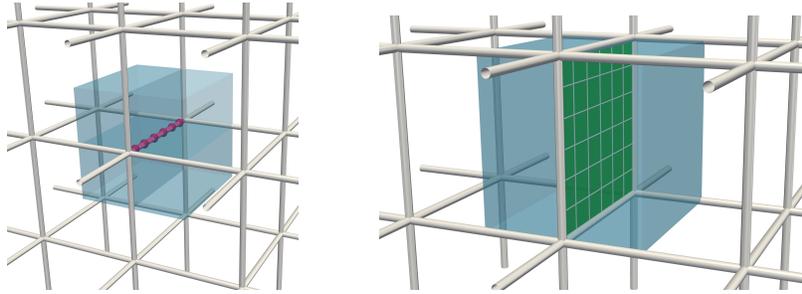


FIG. 5. Three-dimensional version of the slab Ω_e^l corresponding to an edge (left) and a graphical representation of the slab Ω_f^l corresponding to the face f (right).

The slab variant is computationally cheaper and can be proven analogously to the standard version with no modifications. However, as for the variant with local Neumann problems, the coarse space dimension can be larger.

8. Implementation remarks. The classical GDSW coarse space can be implemented algebraically. However, the new coarse space V_{AGDSW} and the variant $V_{\text{AGDSW-S}}$ require the local subdomain stiffness matrices, which cannot be extracted from the global stiffness matrix K . On the other hand, the matrix in the right hand side of the generalized eigenvalue problem (5) can be extracted from the fully assembled stiffness matrix K . Except for the slab variant, the matrix in the left hand side of the generalized eigenvalue problem (5) can be easily computed from the local (non-overlapping) stiffness matrices. For the slab variant, stiffness matrices on slabs need to be assembled. In the variant with local Neumann problems, $V_{\text{AGDSW-S}}$, the implementation is further simplified, since the discrete harmonic extensions are then local to the subdomains and can be computed in parallel. Furthermore, numerical results suggest that GDSW and the adaptive variant only require a simple interface partitioning (components can be disconnected), which facilitates the implementation.

In [28, 27, 26], a parallel implementation of GDSW was considered for various model problems. In a future parallel implementation of AGDSW, we expect the setup of the generalized eigenvalue problems to be the bottleneck. Note that in [24] the inexact solution of the related generalized eigenvalue problems using LobPCG [42] was successful.

9. Numerical results. We present numerical results for the discretized variational problem (1), $f \equiv 1$, and several coefficient functions. Except for the test case in Figure 9 and Table 4, the computational domain is always the unit cube with a zero Dirichlet condition prescribed on its boundary.

We discretize (1) using piecewise trilinear basis functions on voxels or piecewise linear basis functions on tetrahedra and solve the resulting linear system with the preconditioned conjugate gradient (PCG) method and a relative stopping criterion $\|r^{(k)}\|_2 / \|r^{(0)}\|_2 < 10^{-8}$, where $r^{(0)}$ and $r^{(k)}$ are the initial and the k -th unpreconditioned residuals. The reported condition number is the estimate obtained during the last iteration of the PCG method using the Lanczos method [49, ch. 6.7.3].

In the case of voxels, we always consider a cubic domain that is partitioned into smaller cubes. As for tetrahedra, we always partition the domain into subdomains using METIS [34].

We consider the adaptive coarse spaces based on GDSW and the vertex-based and wirebasket coarse spaces by Eikeland, Marcinkowski, and Rahman in [16].

By V_{GDSW} and $V_{\text{AGDSW}}^{\text{tol}_\varepsilon, \text{tol}_\mathcal{F}}$ we denote the GDSW and adaptive GDSW coarse spaces, respectively. Note that $V_{\text{AGDSW}}^{0,0} = V_{\text{GDSW}}$. The variant, which uses a

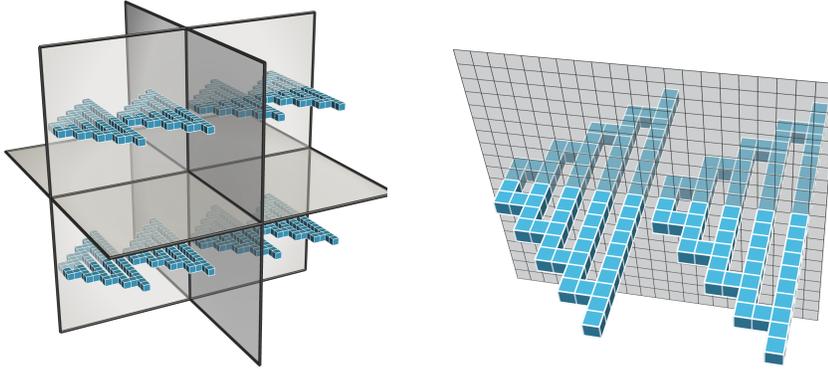


FIG. 6. Discontinuous coefficient function A with coefficients intersecting faces. The blue color corresponds to a coefficient of $A_{\max} = 10^6$ and the remainder is set to $A_{\min} = 1.0$. The left image shows the coefficient function for 2^3 subdomains. One face is highlighted in the right image. The computational domain is a cube discretized using trilinear elements; $1/H = 2$; $H/h = 21$; $\delta = 1h$.

V_0	Coefficient function A from Figure 7				
	\mathcal{F} slab width w	it.	κ	$\dim V_0$	$(\mathcal{V}, \mathcal{E}, \mathcal{F})$
V_{AGDSW}	$1h$	51	30.7	55	(1,6,48)
V_{AGDSW}	$2h$	45	32.7	47	(1,6,40)
V_{AGDSW}	$4h$	45	32.6	39	(1,6,32)
V_{AGDSW}	$6h$	49	32.6	31	(1,6,24)
V_{AGDSW}	$8h$	51	32.6	23	(1,6,16)
V_{AGDSW}	$21h$	51	32.6	23	(1,6,16)
V_{VB}	—	47	30.2	49	(1,0,48)

TABLE 1

Results for the coefficient function in Figure 6: slab width, iteration counts, condition number, and resulting coarse space dimension for different coarse spaces. A tolerance for the selection of the eigenfunctions of 10^{-2} was used for V_{AGDSW} and 10^{-3} for V_{VB} . The full slab width was used for the edge eigenvalue problems of V_{AGDSW} . $1/H = 2$, $H/h = 21$, and $\delta = 1h$; maximum coefficient $A_{\max} = 10^6$; relative stopping criterion $\|r^{(k)}\|_2 / \|r^{(0)}\|_2 < 10^{-8}$.

scaled mass matrix in the right hand side of the eigenvalue problem is denoted by $V_{\text{AGDSW-M}}^{\text{tol}}$; see subsection 7.1. The variant using the sum of local Neumann problems on the left hand side of the eigenvalue problem is denoted by $V_{\text{AGDSW-S}}^{\text{tol}}$; see subsection 7.2. If not mentioned otherwise, the full slab (i.e., the union of subdomains which are adjacent to an edge or face) is used; see also subsection 7.3. By $V_{\text{VB}}^{\text{tol}, \mathcal{E}, \text{tol}, \mathcal{F}}$ and $V_{\text{WB}}^{\text{tol}, \mathcal{E}, \text{tol}, \mathcal{F}}$ we denote the vertex-based and wirebasket coarse spaces from [16].

We begin by showing results for V_{AGDSW} and two coefficient functions by varying the width of the slab in order to highlight the effect of the harmonic extensions in the generalized eigenvalue problems of adaptive GDSW; the same behavior can be observed for the OS-ACMS coarse space introduced in [24]. We then show results for some realistic coefficient functions and, finally, some averaged results for random coefficient functions.

9.1. Varying slab widths. In this section, we investigate the effect of varying slab widths for V_{AGDSW} ; cf. subsection 7.3. Instead of employing the discrete harmonic extensions in the eigenvalue problems on the union of the adjacent subdomains (of an edge or face), we restrict it to a slab around the edge or face. An advantage of using small slabs is the reduced computational cost of computing the discrete harmonic extension. However, we then weaken the detection of connected high coefficient components. The smaller the slab, the fewer connections can be detected.

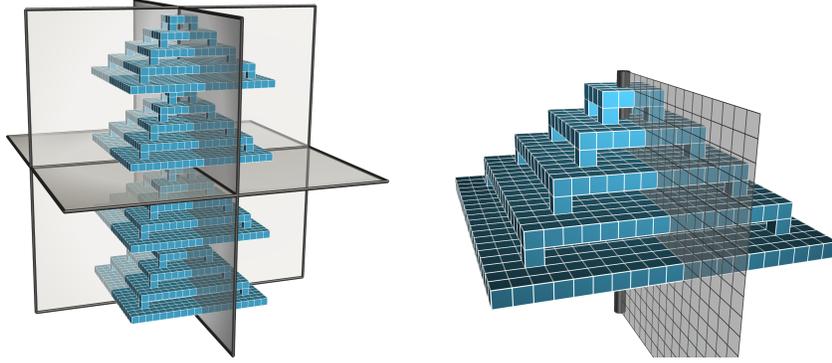


FIG. 7. Discontinuous coefficient function A with coefficients intersecting edges and faces. The blue color corresponds to a coefficient of $A_{\max} = 10^6$ and the remainder is set to $A_{\min} = 1.0$. The left image shows the coefficient function for 2^3 subdomains; each vertical edge is intersected by two high-coefficient connected components. One such component, its corresponding edge, and one adjacent face are highlighted in the right image. The computational domain is a cube discretized using trilinear elements; $1/H = 2$; $H/h = 21$; $\delta = 1h$.

Coefficient function A from Figure 7					
V_0	\mathcal{E} slab width w	it.	κ	$\dim V_0 (\mathcal{V}, \mathcal{E}, \mathcal{F})$	
V_{AGDSW}	$1h$	71	125.5	45 (1,24,20)	
V_{AGDSW}	$2h$	71	125.5	41 (1,20,20)	
V_{AGDSW}	$4h$	70	125.5	37 (1,16,20)	
V_{AGDSW}	$6h$	70	125.6	33 (1,12,20)	
V_{AGDSW}	$8h$	69	125.7	29 (1, 8,20)	
V_{AGDSW}	$21h$	69	125.7	29 (1, 8,20)	
V_{VB}	—	43	37.9	105 (1,20,84)	

TABLE 2

Results for the coefficient function in Figure 7: slab width, iteration counts, condition number, and resulting coarse space dimension for different coarse spaces. A tolerance for the selection of the eigenfunctions of 10^{-2} was used for V_{AGDSW} and 10^{-3} for V_{VB} . The full slab width was used for the face eigenvalue problems of V_{AGDSW} . $1/H = 2$, $H/h = 21$, and $\delta = 1h$; maximum coefficient $A_{\max} = 10^6$; relative stopping criterion $\|r^{(k)}\|_2 / \|r^{(0)}\|_2 < 10^{-8}$.

Therefore, we consider the coefficient function in Figure 6, where high coefficient components intersect faces. The results for various slab widths are listed in Table 1. As can be seen, an increasing slab width yields a decreasing coarse space dimension, since fewer face eigenfunctions are required. A slab of width $8h$ covering the complete high coefficient component is sufficient to yield the same result as for the maximum width of $21h$ (in which case the slab is equal to the union of the adjacent subdomains).

Next, we consider an example for edges; see Figure 7. The numerical results in Table 2 show that using larger slabs in the edge eigenvalue problem reduced the number of edge eigenfunctions. For a slab width of $1h$ a total of 24 edge functions are included in V_{AGDSW} . This number reduces to a minimum of 8 edge functions for a slab width of $8h$.

We conclude that, for certain coefficient functions, the inclusion of the discrete harmonic extension in the eigenvalue problem can significantly reduce the coarse space dimension.

9.2. Realistic and random coefficient functions. In the following, we consider three coefficient functions which exhibit structures that are more likely to be encountered in realistic applications.

Figure 8 depicts 100 beams of high coefficients intersecting a cube that is partitioned into 125 subdomains. As is evident from Table 3 the classical GDSW method

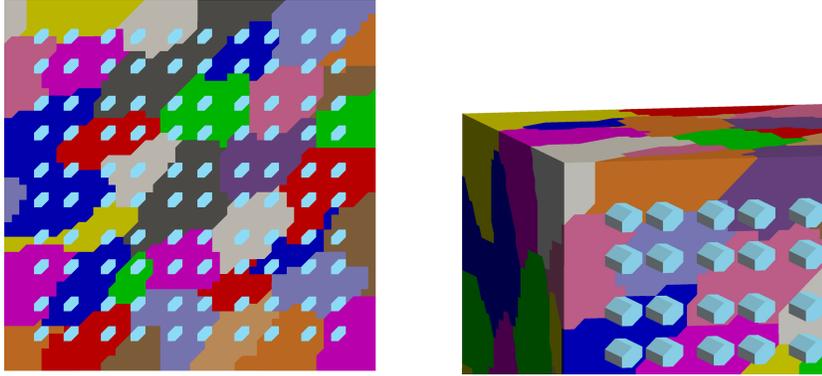


FIG. 8. Cross section (left) of a domain decomposition of a cube and a discontinuous coefficient function A with beams of high coefficients (light blue) crossing the domain. The beams of high coefficients do not touch the domain boundary. The light blue color corresponds to a coefficient of $A_{\max} = 10^6$ and the remainder is set to $A_{\min} = 1.0$. Number of subdomains: 125; number of nodes: 132 651; $\delta = 1h$. Structured tetrahedral mesh; unstructured domain decomposition (METIS).

V_0	Coefficient function A from Figure 8						$\dim V_0 / \text{dof}$
	$\text{tol}_{\mathcal{E}}$	$\text{tol}_{\mathcal{F}}$	it.	κ	$\dim V_0 (\mathcal{V}, \mathcal{E}, \mathcal{F})$		
V_{GDSW}	—	—	1 060	467 954.2	1 987	(564, 790, 633)	1.50%
V_{AGDSW}	10^{-2}	10^{-2}	60	25.1	2 763	(564, 881, 1 318)	2.08%
$V_{\text{AGDSW-S}}$	10^{-2}	10^{-2}	60	25.1	2 763	(564, 881, 1 318)	2.08%
$V_{\text{AGDSW-M}}$	10^{-3}	10^{-3}	60	25.1	2 763	(564, 881, 1 318)	2.08%
V_{VB}	10^{-2}	10^{-2}	58	25.4	3 336	(564, 348, 2 424)	2.51%
V_{WB}	10^{-2}	10^{-2}	50	26.4	5 189	(564, 4 156, 469)	3.91%
Slab width $3h$							
$V_{\text{AGDSW-S}}$	10^{-2}	10^{-2}	60	25.4	2 764	(564, 881, 1 319)	2.08%

TABLE 3

Results for the coefficient function in Figure 8: iteration counts, condition number, and resulting coarse space dimension for different coarse spaces. Number of subdomains: 125; number of nodes: 132 651; $\delta = 1h$; maximum coefficient $A_{\max} = 10^6$; relative stopping criterion $\|r^{(k)}\|_2 / \|r^{(0)}\|_2 < 10^{-8}$. Structured tetrahedral mesh; unstructured domain decomposition (METIS).

requires 1 060 iterations to converge, while all adaptive coarse spaces converge in 60 iterations or less. In this particular example, the adaptive variants of GDSW lead to an increase of 39.1% in the coarse space dimension with respect to GDSW, whereas V_{VB} has a 67.9% and V_{WB} a 161.1% larger coarse space. Reducing the slab width of $V_{\text{AGDSW-S}}$ to only $3h$ is sufficient to obtain almost identical results.

As a second example, we consider the coefficient function in Figure 9 with several layers of varying coefficients. We note that most of the domain is surrounded by a homogeneous Neumann boundary condition and the Dirichlet boundary does not touch a high coefficient layer; see Figure 9 (center). Despite a condition number of $3.8 \cdot 10^6$, the classical GDSW method requires only 125 iterations to converge, due to the relatively low number of coefficient jumps. For adaptive GDSW, we observe an increase in the coarse space dimension of only 13.7% compared to classical GDSW, while the dimension of V_{VB} is 59.3% larger and the wirebasket coarse space's dimension is more than twice as large as that of classical GDSW; cf. Table 4.

As a final realistic example, we consider a foam-like structure of high coefficients embedded in a cube; cf. Figure 10. We note that the foam structure consists of several disconnected smaller foam structures. The numerical results in Table 5 show that V_{AGDSW} converges within 61 iterations for a tolerance of 0.01, while increasing the coarse space dimension by only 22.5% compared to V_{GDSW} . Using

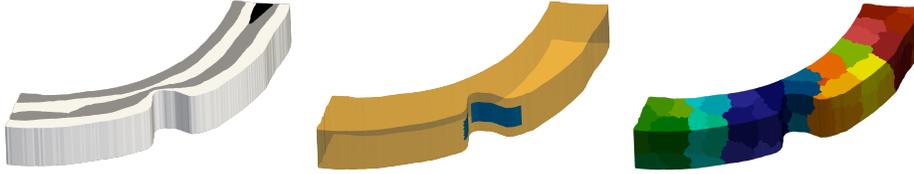


FIG. 9. (left) Discontinuous coefficient function A with coefficient layers of $A = 10^6$ in light gray and an inclusion at the top right with $A = 10^9$ in dark grey. The remainder of the coefficient in white is set to $A_{\min} = 1.0$. (center) Boundary partition for Dirichlet (blue) and Neumann (orange) boundary. (right) Domain decomposition of 50 subdomains. Number of nodes: 56 053; average number of nodes per subdomain: 1313.0; $\delta = 1h$. Unstructured tetrahedral mesh; unstructured domain decomposition (METIS).

V_0	Coefficient function A from Figure 9						dim V_0 ($\mathcal{V}, \mathcal{E}, \mathcal{F}$)	dim V_0 /dof
	$tol_{\mathcal{E}}$	$tol_{\mathcal{F}}$	it.	κ				
V_{GDSW}	—	—	125	3 770 557.2	445	(105, 168, 172)	0.79%	
V_{AGDSW}	10^{-2}	10^{-2}	50	20.1	506	(105, 186, 215)	0.90%	
$V_{\text{AGDSW-S}}$	10^{-2}	10^{-2}	50	20.1	506	(105, 186, 215)	0.90%	
$V_{\text{AGDSW-M}}$	10^{-3}	10^{-3}	50	20.1	506	(105, 186, 215)	0.90%	
V_{VB}	10^{-3}	10^{-3}	49	18.8	709	(105, 41, 563)	1.26%	
V_{WB}	10^{-3}	10^{-3}	41	15.0	964	(105, 854, 5)	1.72%	
Slab width $3h$								
$V_{\text{AGDSW-S}}$	10^{-2}	10^{-2}	50	20.1	506	(105, 186, 215)	0.90%	

TABLE 4

Results for the coefficient function in Figure 9: iteration counts, condition number, and resulting coarse space dimension for different coarse spaces. Number of subdomains: 50; number of nodes: 56 053; average number of nodes per subdomain: 1 313.0; $\delta = 1h$; maximum coefficient $A_{\max} = 10^9$; relative stopping criterion $\|r^{(k)}\|_2 / \|r^{(0)}\|_2 < 10^{-8}$. Unstructured tetrahedral mesh; unstructured domain decomposition (METIS).

a slab width of $3h$ and $V_{\text{AGDSW-S}}$ (and a tolerance of 0.01) results in an increase of the coarse space dimension with respect to GDSW by 23.6% and convergence is achieved within 60 iterations. In contrast, the coarse space V_{VB} leads to an increase in the coarse space dimension of 104.1% while requiring 58 iterations to converge (tolerance: 10^{-5}).

Finally, we present averaged results for 100 randomly generated coefficient functions with an average of 11.1% elements with high coefficients $A_{\max} = 10^6$ (the remainder is set to $A_{\max} = 1$). The results in Table 6 show that, also for random coefficient functions, the adaptive GDSW variants perform well. The largest coarse space dimension 13 665 of an adaptive GDSW variant is attained by restricting $V_{\text{AGDSW-S}}$ to a slab of width $3h$ ($tol_{\mathcal{E}} = tol_{\mathcal{F}} = 0.1$). This amounts to an increase of 36.2% compared to V_{GDSW} . Simultaneously, the largest number of iterations for this setting is 80, whereas classical GDSW did not converge within 2 000 iterations.

10. Conclusion. We have presented a new adaptive coarse space for the overlapping Schwarz method and proved a condition number bound. This bound depends on user prescribed tolerances but is independent of the mesh parameters h , H , and of heterogeneities in the coefficient function A .

At its core AGDSW uses generalized eigenvalue problems on edges and faces which are thus of moderate size compared to some competing approaches. AGDSW always contains the GDSW coarse space and only requires local nonoverlapping stiffness matrices to set up the eigenvalue problems; all other information can be obtained algebraically. Several variants have been presented among which the one in subsection 7.2 facilitates the implementation and reduces the computational complexity by increasing sparsity.

The results in section 9 support the theory and show that using the discrete

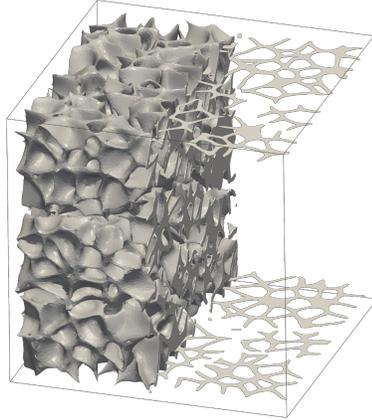


FIG. 10. Partial visualization of an unstructured tetrahedral mesh consisting of several disconnected components of foam-like structures. In the corresponding mesh of a cube, the foam corresponds to a high coefficient of $A_{\max} = 10^6$ and is to $A_{\min} = 1.0$ elsewhere. The high coefficient does not touch the domain boundary. Number of subdomains: 100; number of nodes: 588 958; average number of nodes per subdomain: 6 656.4; $\delta = 1h$. Unstructured tetrahedral mesh; unstructured domain decomposition (METIS).

V_0	Coefficient function A from Figure 10							
	$tol_{\mathcal{E}}$	$tol_{\mathcal{F}}$	it.	κ	$\dim V_0 (\mathcal{V}, \mathcal{E}, \mathcal{F})$			$\dim V_0 / \text{dof}$
V_{GDSW}	—	—	565	1 337 890.8	1 601	(404, 688, 509)	0.27%	
V_{AGDSW}	10^{-2}	10^{-2}	61	30.7	1 962	(404, 741, 817)	0.33%	
V_{AGDSW}	10^{-1}	10^{-1}	54	22.5	2 055	(404, 741, 910)	0.35%	
$V_{\text{AGDSW-S}}$	10^{-2}	10^{-2}	61	30.7	1 963	(404, 742, 817)	0.33%	
$V_{\text{AGDSW-S}}$	10^{-1}	10^{-1}	54	22.7	2 060	(404, 742, 914)	0.35%	
$V_{\text{AGDSW-M}}$	10^{-3}	10^{-3}	61	30.7	1 962	(404, 741, 817)	0.33%	
$V_{\text{AGDSW-M}}$	10^{-2}	10^{-2}	60	30.1	1 966	(404, 741, 821)	0.33%	
V_{VB}	10^{-6}	10^{-6}	551	35 709.5	2 740	(404, 453, 1 883)	0.47%	
V_{VB}	10^{-5}	10^{-5}	58	27.0	3 268	(404, 453, 2 411)	0.55%	
V_{WB}	10^{-6}	10^{-6}	317	11 644.0	5 941	(404, 5 501, 36)	1.01%	
V_{WB}	10^{-5}	10^{-5}	46	19.1	6 195	(404, 5 501, 290)	1.05%	
Slab width $3h$								
$V_{\text{AGDSW-S}}$	10^{-2}	10^{-2}	60	29.7	1 979	(404, 744, 831)	0.34%	
$V_{\text{AGDSW-S}}$	10^{-1}	10^{-1}	48	17.9	2 241	(404, 748, 1 089)	0.38%	

TABLE 5

Results for the coefficient function in Figure 10: iteration counts, condition number, and resulting coarse space dimension for different coarse spaces. Number of subdomains: 100; number of nodes: 588 958; average number of nodes per subdomain: 6 656.4; $\delta = 1h$; maximum coefficient $A_{\max} = 10^6$; relative stopping criterion $\|r^{(k)}\|_2 / \|r^{(0)}\|_2 < 10^{-8}$. Unstructured tetrahedral mesh; unstructured domain decomposition (METIS).

harmonic extension inside the eigenvalue problem (5) can help to reduce the dimension of the coarse space by detecting connected components of high coefficients. Furthermore, we have demonstrated the robustness of AGDSW for various realistic coefficient distributions.

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V_0	Random coefficient function A					
	$tol_{\mathcal{E}}$	$tol_{\mathcal{F}}$	it.	κ	$\dim V_0$	$\dim V_0/\text{dof}$
V_{GDSW}	–	–	>2000 (–)	$2.7 \cdot 10^5$ ($3.7 \cdot 10^5$)	10031.0 (10031)	2.2% (2.2%)
V_{AGDSW}	10^{-1}	10^{-1}	79.4 (88)	52.3 (74.6)	13244.3 (13397)	2.9% (3.0%)
V_{AGDSW}	10^{-2}	10^{-2}	134.6 (158)	166.9 (270.4)	12749.8 (12963)	2.8% (2.9%)
$V_{\text{AGDSW-S}}$	10^{-1}	10^{-1}	74.3 (81)	45.6 (68.8)	13438.0 (13585)	3.0% (3.0%)
$V_{\text{AGDSW-S}}$	10^{-2}	10^{-2}	133.3 (157)	164.5 (270.4)	12838.9 (13043)	2.8% (2.9%)
$V_{\text{AGDSW-M}}$	10^{-2}	10^{-2}	113.1 (131)	114.7 (169.5)	12825.2 (13007)	2.8% (2.9%)
$V_{\text{AGDSW-M}}$	10^{-3}	10^{-3}	140.6 (165)	187.5 (311.2)	12674.0 (12879)	2.8% (2.9%)
V_{VB}	10^{-3}	10^{-3}	74.5 (87)	54.6 (132.8)	21576.2 (22147)	4.8% (4.9%)
V_{VB}	10^{-5}	10^{-5}	74.8 (87)	55.6 (132.8)	21554.4 (22144)	4.8% (4.9%)
V_{WB}	10^{-3}	10^{-3}	60.2 (72)	36.0 (82.8)	27659.4 (27884)	6.1% (6.2%)
V_{WB}	10^{-5}	10^{-5}	61.2 (72)	39.5 (100.2)	27645.8 (27880)	6.1% (6.2%)
Slab width $3h$						
$V_{\text{AGDSW-S}}$	10^{-1}	10^{-1}	71.7 (80)	42.2 (62.2)	13526.3 (13665)	3.0% (3.0%)
$V_{\text{AGDSW-S}}$	10^{-2}	10^{-2}	122.6 (137)	137.4 (270.4)	12924.5 (13109)	2.9% (2.9%)

TABLE 6

Averaged results for 100 random coefficient functions (average high coefficient density: 11.08%): tolerance for the selection of the eigenfunctions, iteration counts, condition number, and resulting coarse space dimension for different coarse spaces; maximum in brackets. Number of subdomains: 512; number of nodes: 452522; average number of nodes per subdomain: 1174.4; $\delta = 1h$; maximum coefficient $A_{\max} = 10^6$; relative stopping criterion $\|r^{(k)}\|_2/\|r^{(0)}\|_2 < 10^{-8}$. Unstructured tetrahedral mesh; unstructured domain decomposition (METIS). V_{GDSW} never converged within 2000 iterations.

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